Solution of systems of algebraic equations in CFD

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The Poisson equation must be solved in every time step

in Ψ - ω method:

$$\Delta \psi = -\omega \longrightarrow \psi$$

in pressure based methods:

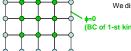
$$\Delta P = \nabla \cdot f \longrightarrow P$$

A simple 2D example

The computational domain:

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = Q$$

We discretize this by using compass notations:



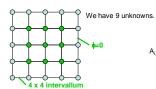
$$\frac{1}{\varDelta x} \left(\frac{\phi_E - \phi_P}{\varDelta x} - \frac{\phi_P - \phi_W}{\varDelta x} \right) + \frac{1}{\varDelta y} \left(\frac{\phi_N - \phi_P}{\varDelta y} - \frac{\phi_P - \phi_S}{\varDelta y} \right) = Q_P$$

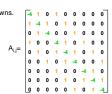
$$\Delta x = \Delta y = h$$

On isotropic mesh:
$$\Delta x = \Delta y = h$$
 $\phi_S + \phi_W - 4\phi_P + \phi_E + \phi_N = h^2 Q_P$

In matrix form

$$\phi_S + \phi_W - 4\phi_P + \phi_E + \phi_N = h^2 Q_P$$





The system now reads:

 $A_{i,j} \varphi_i = Q_i$

The number of unknowns for 101x101 mesh N=104, therefore the number of elements of matrix A is 108.

Gauss elimination

As efficient as any other method for a general case, but it does not make use of the favorable characteristics of the matrix.

 $A_{2,1}$ $A_{2,2}$ $A_{2,3}$ $\begin{pmatrix} A_{3,1} & A_{3,2} & A_{3,3} \end{pmatrix}$ second row up to the column N-1.

1-st step Elimination: A_{21}/A_{11} times the first row is subtracted from the second row.

Repeated for $\left(egin{array}{ccc} A_{1,1} & A_{1,2} & A_{1,3} \end{array}
ight)$ The first element of the second row every further will become 0. Similarly, we eliminate the other elements of the

2-nd step Backsubstitution:

$$\begin{pmatrix} U_{1,1} & U_{1,2} & U_{1,3} \\ 0 & U_{2,2} & U_{2,3} \\ 0 & 0 & U_{3,3} \end{pmatrix}$$

$$\phi_n = \frac{Q_n}{U_{nn}}$$

$$\phi_i = \frac{Q_i - \sum_{k=k+1}^{N} U_{k,i} \phi_k}{U_{i:i}}$$

The operation cost of the method is N3/3, out of which the back substitution requires only N²/2 operations. Even if A is sparse U is not spares. The total memory requirement on a 2D mesh of 101x101 nodes is 400 Mb. We don't need such an accurate solution because the discretization error is large anyway.

Iterative methods

The solution is refined step by step: approximation of ϕ in the n-th step is ϕ^n .

By omitting the vector notations: $A\phi^n = Q - \rho^n$ ρ^n : residual

 $A \varepsilon^n = A(\phi - \phi^n) = Q - (Q - \rho^n) = \rho^n$

Thus, the exacts solution is obtained in n=1 step if A matrix is solved for the error. We can

 $M \phi^{n+1} = N\phi^n + Q$

For the converged solution: $\phi^{n+1} = \phi^n = \phi$, therefore: A = M - N

Let's subtract $M\phi^{\,n}\,$ from both sides:

$$M(\phi^{n+1} - \phi^n) = N \phi^n + Q - M \phi^n = Q - A \phi^n = \rho^n$$

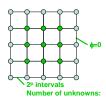
 $M \delta^n = \rho^n$ This is the correction equation. correction: δ^n

The better M approximates A is the faster the method converges. M must be easy to solve eg. diagonal, tri-diagonal, or a Δ matrix.

Jacobi iteration

$$\phi_S^n + \phi_W^n - 4\phi_P^{n+1} + \phi_E^n + \phi_N^n = h^2 Q_P \qquad M \text{ is a diagonal matrix}.$$

$$\phi_P^{n+1} = \frac{1}{4} \left(\phi_S^n + \phi_W^n + \phi_E^n + \phi_N^n - h^2 Q_P \right)$$



(2p-1)2.

Example program:

- The program...
- Major characteristics of the result...
 Required number of iterations ...

Gauss-Seidel relaxation



These terms are already known due to the calculation

$$\phi_P^{n+1} = \frac{1}{4} \left(\phi_S^n + \phi_W^{n+1} + \phi_E^n + \phi_N^{n+1} - h^2 Q_P \right)$$

- It requires halve as much iterations ...
- and halve as much memory.
- The error is asymmetrically distributed.

Line relaxation



$$\phi_{S}^{n} + \phi_{W}^{n+1} - 4\phi_{P}^{n+1} + \phi_{E}^{n+1} + \phi_{N}^{n+1} = h^{2}Q_{P}$$
these values

these obtained line by line from the tri-diagonal system solved by the Thomas algorithm.

Note - Much more efficient methods based on the tri-diagonal solver also exist: Operator Splitting (or Alternating Direction Implicit, ADI) methods.

The problem:

The above mentioned methods are only smoothing the solution.

The boundary effects need a very long time to penetrate the computational domain.

We need to use coarser meshes too. The first estimates of the correction can be obtained on a coarser mesh, than can be refined on the fine mesh.

Multigrid method

The correction equation for a simplified 1D problem:

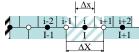
$$\frac{\partial^2 \phi}{\partial x^2} = Q$$

$$\frac{1}{\Delta x^2} (\phi_{i-1} - 2\phi_i + \phi_{i+1}) = Q_i$$

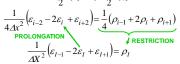
$$\frac{1}{\Delta x^2} (\phi_{i-1}^n - 2\phi_i^n + \phi_{i+1}^n) = Q_i - \rho_i^n$$

$$\frac{1}{4r^2} \left(\varepsilon_{i-1}^n - 2\varepsilon_i^n + \varepsilon_{i+1}^n \right) = \rho_i^n$$

We omit the iteration indices: $\frac{1}{4r^2} \left(\varepsilon_{i-1} - 2\varepsilon_i + \varepsilon_{i+1} \right) = \rho_i$



$$\begin{split} \frac{1}{4x^2} & \left(\frac{1}{2} \varepsilon_{i-2} - \varepsilon_{i-1} + \frac{1}{2} \varepsilon_i + \varepsilon_{i-1} - 2 \varepsilon_i + \varepsilon_{i+1} + \frac{1}{2} \varepsilon_i - \varepsilon_{i+1} + \frac{1}{2} \varepsilon_{i+2} \right) = \\ \text{these terms are cancelled} & = \frac{1}{2} \rho_{i-1} + \rho_i + \frac{1}{2} \rho_{i+1} \end{split}$$



Generalization to 2D or 3D:

Restriction:



- 1. Restriction: $\rho_i \to \rho_i$ 2. Calculation of ϵ_i . Eg. in 3D we have an 8 fold reduced number of unknowns. 3. Prolongation of ϵ_i to the fine mesh. (ϵ_i) ,
- 4. Smoothing on the fine mesh.

Why shouldn't we use an even more coarse mesh when calculating ε_{l} ?

- 1. Evaluation of the residuals on the finest mesh.
- Consecutive restrictions of ρ to every coarser mesh.
- Solution of the system on the coarsest mesh. (Even by using a direct method.)
- Consecutively for every finer mesh:
- Prolongation of ε
- Smoothing (Eg. by using Gauss-Seidel relaxation.)
- Correction of φ (Only on the finest mesh).

	С	ompu	ıtatior	nal co	st	
Numb	er of iteration	ns in 2D:				
р	N _{line}	N	Jacobi	G-S	Line rlx.	Multigr
3	7	49	40	20	10	11
4	15	225	160	80	40	24
5	31	961	640	320	160	38
6	63	3969	2560	1280	640	44
7	127	16129	10240	5120	2560	46
Numb	er of operation	ons / N:				
р	N _{line}	N	Jacobi	G-S	Line rlx.	Multigr
3	7	49	200	100	50	220
4	15	225	800	400	200	480
5	31	961	3200	1600	800	760
6	63	3969	12800	6400	3200	880
7	127	16129	51200	25600	12800	920
				On fine mes	hes multigri	d prevail