## Solution of systems of algebraic equations in CFD

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## In matrix form

$$
\phi_{S}+\phi_{V}-4 \phi_{P}+\phi_{E}+\phi_{N}=h^{2} Q_{P}
$$



The system now reads:

$$
A_{i, j} \varphi_{i}=Q_{i}
$$

The number of unknowns for $101 \times 101$ mesh $\mathrm{N}=10^{4}$, therefore the number of elements of matrix A is $10^{8}$.

The Poisson equation must be solved in every time step
in $\Psi-\omega$ method:
$\Delta \psi=-\omega \longrightarrow \psi$
in pressure based methods:

## Gauss elimination

As efficient as any other method for a general case, but it does not make use of the favorable characteristics of the matrix.
1-st step Elimination: $\quad A_{21} / A_{11}$ times the first row is
\(\left.\left($$
\begin{array}{lll}A_{1,1} & A_{1,2} & A_{1,3} \\
A_{2,1} & A_{2,2} & A_{2,3} \\
A_{3,1} & A_{3,2} & A_{3,3}\end{array}
$$\right) \begin{array}{l}subtracted from the second row. <br>
The first element of the second row <br>
will become 0. Similarly, we <br>
eliminate the other elements of the <br>

second row up to the column N-1.\end{array}\right\}\)| Repeated for |
| :--- |
| every further |
| rows. |

$$
\begin{aligned}
& \text { 2-nd step Backsubstitution: } \\
& \left(\begin{array}{ccc}
U_{1,1} & U_{1,2} & U_{1,3} \\
0 & U_{2,2} & U_{2,3} \\
0 & 0 & U_{3,3}
\end{array}\right)
\end{aligned} \phi_{n}=\frac{Q_{n}}{U_{n n}} .
$$

The operation cost of the method is $\mathrm{N}^{3} / 3$, out of which the back substitution requires only $\mathrm{N}^{2} / 2$ operations. Even if $A$ is sparse $U$ is not spares. The total memory requirement on a 2D mesh of $101 \times 101$ nodes is 400 Mb . We don't need such an accurate solution because the discretization error is large anyway.

| A simple 2D example $\text { The computational domain: } \frac{\partial^{2} \phi}{\partial x^{2}}+\frac{\partial^{2} \phi}{\partial y^{2}}=Q$ <br> We discretize this by using compass notations: $\frac{1}{\Delta x}\left(\frac{\phi_{E}-\phi_{P}}{\Delta x}-\frac{\phi_{P}-\phi_{W}}{\Delta x}\right)+\frac{1}{\Delta y}\left(\frac{\phi_{N}-\phi_{P}}{\Delta y}-\frac{\phi_{P}-\phi_{S}}{\Delta y}\right)=Q_{P}$ <br> On isotropic mesh: $\Delta x=\Delta y=h \quad \phi_{S}+\phi_{W}-4 \phi_{P}+\phi_{E}+\phi_{N}=h^{2} Q_{P}$ |
| :---: |
|  |  |
|  |  |

## Iterative methods

The solution is refined step by step: approximation of $\phi$ in the $n$-th step is $\phi^{n}$. By omitting the vector notations: $A \phi^{n}=Q-\rho^{n} \quad \rho^{n}$ : residual
The error: $\quad \varepsilon^{n}=\phi-\phi^{n}$
$A \varepsilon^{n}=A\left(\phi-\phi^{n}\right)=Q-\left(Q-\rho^{n}\right)=\rho^{n} \quad \begin{aligned} & \text { Thus, the same matrix describes } \\ & \text { the error }\end{aligned}$ the error

Iterative methods:

$$
M \phi^{n+1}=N \phi^{n}+Q
$$

For the converged solution: $\phi^{n+1}=\phi^{n}=\phi \quad$, therefore: $A=M-N$
Let's subtract $M \phi^{n}$ from both sides:

$$
M\left(\phi^{n+1}-\phi^{n}\right)=N \phi^{n}+Q-M \phi^{n}=Q-A \phi^{n}=\rho^{n}
$$

## Jacobi iteration

$\phi_{S}^{n}+\phi_{W}^{n}-4 \phi_{P}^{n+1}+\phi_{E}^{n}+\phi_{N}^{n}=h^{2} Q_{P} \quad M$ is a diagonal matrix.
$\phi_{P}^{n+1}=\frac{1}{4}\left(\phi_{S}^{n}+\phi_{W}^{n}+\phi_{E}^{n}+\phi_{N}^{n}-h^{2} Q_{P}\right)$


## Example program

The program..
Major characteristics of the result.. Required number of iterations ...

## Gauss-Seidel relaxation



## Multigrid method

The correction equation for a simplified 1D problem:

$$
\begin{gathered}
\frac{\partial^{2} \phi}{\partial x^{2}}=Q \\
\frac{1}{\Delta x^{2}}\left(\phi_{i-1}-2 \phi_{i}+\phi_{i+1}\right)=Q_{i} \\
\frac{1}{\Delta x^{2}}\left(\phi_{i-1}^{n}-2 \phi_{i}^{n}+\phi_{i+1}^{n}\right)=Q_{i}-\rho_{i}^{n} \\
\frac{1}{\Delta x^{2}}\left(\varepsilon_{i-1}^{n}-2 \varepsilon_{i}^{n}+\varepsilon_{i+1}^{n}\right)=\rho_{i}^{n}
\end{gathered}
$$

We omit the iteration indices: $\frac{1}{\Delta x^{2}}\left(\varepsilon_{i-1}-2 \varepsilon_{i}+\varepsilon_{i+1}\right)=\rho_{i}$

$\frac{1}{\Delta x^{2}}\left(\frac{1}{2} \varepsilon_{i-2}-\varepsilon_{i-1}+\frac{1}{2} \varepsilon_{i}+\underline{\varepsilon_{i-1}}-2 \varepsilon_{i}+\underline{\varepsilon_{i+1}}+\frac{1}{2} \varepsilon_{i} \underline{-\varepsilon_{i+1}}+\frac{1}{2} \varepsilon_{i+2}\right)=$
these terms are cancelled $=\frac{1}{2} \rho_{i-1}+\rho_{i}+\frac{1}{2} \rho_{i+1}$

$$
\begin{aligned}
& \frac{1}{4 \Delta x^{2}}\left(\varepsilon_{i-2}-2 \varepsilon_{i}+\varepsilon_{i+2}\right)=\underbrace{\frac{1}{4}\left(\rho_{i-1}+2 \rho_{i}+\rho_{i+1}\right)}_{\substack{\text { PROLONGATION }}} \\
& \frac{1}{\Delta X^{2}}\left(\varepsilon_{I-1}-2 \varepsilon_{I}+\varepsilon_{I+1}\right)=\rho_{I}
\end{aligned}
$$

## Line relaxation



## Generalization to 2D or 3D:

Restriction:


1. Restriction: $\rho_{i} \rightarrow \rho_{t}$
2. Calculation of $\varepsilon_{1}$. Eg. in 3D we have an 8 fold reduced number of unknowns.
3. Prolongation of $\varepsilon_{\text {t }}$ to the fine mesh. $\left(\varepsilon_{i}\right)$,
4. Smoothing on the fine mesh.

Why shouldn't we use an even more coarse mesh when calculating $\varepsilon_{1}$ ?
Operator Splitting (or Alternating Direction Implicit, ADI) methods.

1. Evaluation of the residuals on the finest mesh.
2. Consecutive restrictions of $\rho$ to every coarser mesh.

The problem:
The above mentioned methods are only smoothing the solution.
The boundary effects need a very long time to penetrate the computational domain
3. Solution of the system on the coarsest mesh. (Even by using a direct method.)
4. Consecutively for every finer mesh:

Prolongation of $\varepsilon$

- Smoothing (Eg. by using Gauss-Seidel relaxation.)

We need to use coarser meshes too. The first estimates of the correction can be obtained on a coarser mesh, than can be refined on the fine mesh.

| Computational cost |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Number of iterations in 2D: |  |  |  |  |  |  |
| p | $\mathrm{N}_{\text {line }}$ | N | Jacobi | G-S | Line rlx. | Multigrid |
| 3 | 7 | 49 | 40 | 20 | 10 | 13 |
| 4 | 15 | 225 | 160 | 80 | 40 | 33 |
| 5 | 31 | 961 | 640 | 320 | 160 | 59 |
| 6 | 63 | 3969 | 2560 | 1280 | 640 | 75 |
| 7 | 127 | 16129 | 10240 | 5120 | 2560 | 79 |
| Müveletigény / N : |  |  |  |  |  |  |
| p | $\mathrm{N}_{\text {line }}$ | N | Jacobi | G-S | Line rlx. | Multigrid |
| 3 | 7 | 49 | 200 | 100 | 50 | 260 |
| 4 | 15 | 225 | 800 | 400 | 200 | 660 |
| 5 | 31 | 961 | 3200 | 1600 | 800 | 1180 |
| 6 | 63 | 3969 | 12800 | 6400 | 3200 | 1500 |
| 7 | 127 | 16129 | 51200 | 25600 | 12800 | 1580 |
| On fine meshes multigrid prevails! |  |  |  |  |  |  |

