

Turbulence  
II.

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Balogh

Scales

TKE eq.

Modelling

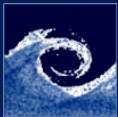
Boundaries  
Inlet

# Turbulence modelling II.

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November 13, 2018



# Many scales of turbulence

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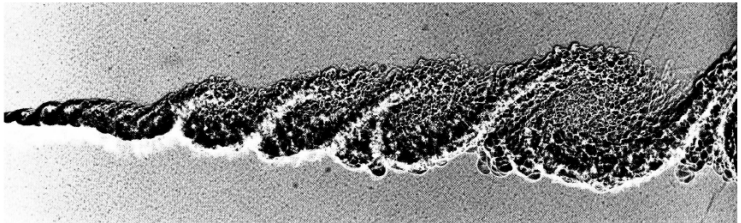
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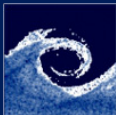
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Density variation visualise the different scales of turbulence in a mixing layer



Goal: Try to find some rules about the properties of turbulence at different scales



# Kinetic energy

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Kinetic energy:

$$E \stackrel{\text{def}}{=} \frac{1}{2} u_i u_i \quad (1)$$

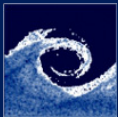
Its Reynolds decomposition:

$$E = \frac{1}{2} u_i u_i = \frac{1}{2} (\overline{u_i u_i} + 2u'_i \overline{u_i} + u'_i u'_i) \quad (2)$$

Its Reynolds average

$$\overline{E} = \underbrace{\frac{1}{2} (\overline{u_i u_i})}_{\hat{E}} + \underbrace{\frac{1}{2} (\overline{u'_i u'_i})}_k = \hat{E} + k \quad (3)$$

- The kinetic energy of the mean flow:  $\hat{E}$
- The kinetic energy of the turbulence:  $k$  (Turbulent Kinetic Energy, TKE)



# Richardson energy cascade

## The poem

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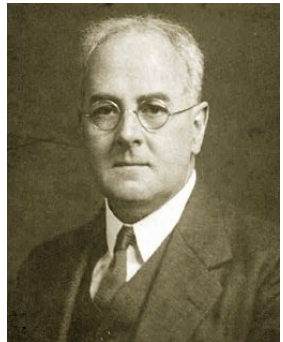
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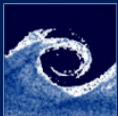
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Lewis Fry Richardson (1920):

*"Big whirls have little whirls,  
that feed on their velocity;  
and little whirls have lesser whirls,  
and so on to viscosity."*

*"Nagy örvény kisebbet plántál,  
melyet sebességével táplál;  
majd az még kisebbet szülén,  
viszkozitásba tűnik szürkén."*





# Richardson energy cascade

## Vortex scales

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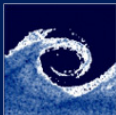
## High Re flow is investigated

- Typical velocity of the flow  $\mathcal{U}$
- Typical length scale of the flow  $\mathcal{L}$
- Corresponding Reynolds number ( $\mathcal{Re} = \frac{\mathcal{U}\mathcal{L}}{\nu}$ ) is high

## Turbulence is made of vortices of different sizes

Each class of vortex has:

- length scale:  $l$
- velocity scale:  $u(l)$
- time scale:  $\tau(l) = l/u(l)$



# Richardson energy cascade

## The big scales

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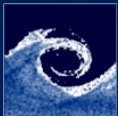
### Biggest vortices

- size  $l_0 \sim \mathcal{L}$
- velocity  $u_0 = u_0(l_0) \sim u' = \sqrt{2/3k} \sim \mathcal{U}$

$\Rightarrow Re = \frac{u_0 l_0}{\nu}$  is also high

### Fragmentation of the big vortices

- High  $Re$  corresponds to low viscous stabilisation
- Big vortices are unstable
- Big vortices break up into smaller ones



# Richardson energy cascade

## To the small scales

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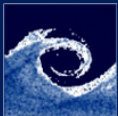
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### Inertial cascade

- As long as  $Re(l)$  is high, inertial forces dominate, the break up continues
- At small scales  $Re(l) \sim 1$  viscosity starts to be important
  - The kinetic energy of the vortices dissipates into heat



# Richardson energy cascade

Connection between small and large scales

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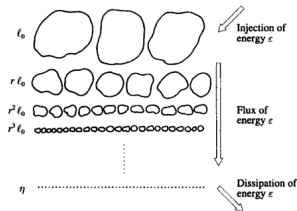
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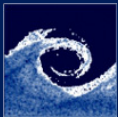
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## Dissipation equals production

- Dissipation is denoted by  $\varepsilon$
- Because of the cascade can be characterised by large scale motion
- Dissipation:  $\varepsilon \sim \frac{\text{kin. energy}}{\text{timescale}}$  at the large scales
  - By formula:  $\varepsilon = \frac{u_0^2}{l_0/u_0} = \frac{u_0^3}{l_0}$







# Transport equation of $k$

## Definitions 1

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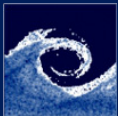
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### NS symbol

For the description of development rules, it is useful to define the following NS symbol:

$$NS(u_i) \stackrel{\text{def}}{=} \partial_t u_i + u_j \partial_j u_i = - \underbrace{\frac{1}{\rho} \partial_i p + \nu \partial_j s_{ij}}_{\partial_j t_{ij}} \quad (4)$$

where:  $s_{ij} \stackrel{\text{def}}{=} \frac{1}{2}(\partial_i u_j + \partial_j u_i)$  is the deformation (rate of strain) part of the derivative tensor  $\partial_j u_i$ .



# Transport equation of $k$

## Definitions 2

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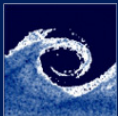
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Let us repeat the development of the Reynolds equation!

$$\overline{NS(\overline{u_i} + u'_i)} \quad (5)$$

$$\partial_t \overline{u_i} + \overline{u_j} \partial_j \overline{u_i} = \partial_j \underbrace{\left[ -\frac{1}{\rho} \overline{p} \delta_{ij} + \nu \overline{s}_{ij} - \overline{u'_i u'_j} \right]}_{\overline{T_{ij}}} \quad (6)$$



# The TKE equation

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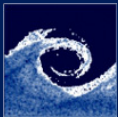
Taking the trace of

$$\overline{(NS(u_i) - \overline{NS(u_i)})u'_j} + \overline{(NS(u_j) - \overline{NS(u_j)})u'_i}$$

$$\partial_t k + \overline{u_j} \partial_j k = \underbrace{-a_{ij} \overline{s'_{ij}}}_{\text{Production}} + \underbrace{\partial_j \left[ u'_j \left( \frac{p'}{\rho} + k' \right) - \nu \overline{u'_i s'_{ij}} \right]}_{\text{Transport}} - \underbrace{\varepsilon}_{\text{Dissipation}} \quad (7)$$

- Dissipation:  $\varepsilon \stackrel{\text{def}}{=} 2\nu \overline{s'_{ij} s'_{ij}}$
- Anisotropy tensor:  $a_{ij} \stackrel{\text{def}}{=} \overline{u'_i u'_j} - \frac{1}{3} \underbrace{\overline{u'_l u'_l}}_{2k} \delta_{ij}$

Deviator part of the Reynolds stress tensor



# The TKE equation

## Meaning of the terms

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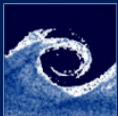
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### Production

- Expression:  $\mathcal{P} \stackrel{\text{def}}{=} -a_{ij} \overline{s_{ij}}$
- Transfer of kinetic energy from mean flow to turbulence
  - The same term with opposite sign in the equation for kin. energy of mean flow
- The mechanism to put energy in the 'Richardson' cascade
- Happens at the large scales



# The TKE equation

## Meaning of the terms (contd.)

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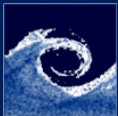
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## Dissipation

- Expression:  $\varepsilon \stackrel{\text{def}}{=} 2\nu \overline{s'_{ij}s'_{ij}}$
- Conversion of kinetic energy of turbulence to heat
  - Work of the viscous stresses at small scale ( $s'_{ij}$ )
- The mechanism to draw energy from the 'Richardson' cascade
- Happens at the small scales

$\mathcal{P} = \varepsilon$  if the turbulence is homogeneous (isotropic), as in the "Richardson" cascade



# The TKE equation

## Meaning of the terms (contd.)

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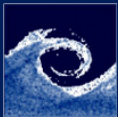
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## Transport

- Expression:  $\partial_j \left[ \overline{u'_j \left( \frac{p'}{\rho} + k' \right)} - \nu \overline{u'_i s'_{ij}} \right]$
- Transport of turbulent kinetic energy in space
  - The expression is in the form of a divergence ( $\partial_j \square_j$ )
  - Divergence can be reformulated to surface fluxes (G-O theorem)



# Idea of RANS modelling

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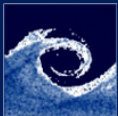
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- Solving the Reynolds averaged NS for the averaged variables  $(\bar{u}, \bar{v}, \bar{w}, \bar{p})$
- The Reynolds stress tensor  $\overline{u'_i u'_j}$  is unknown and has to be modelled
- Modelling should use the available quantities  $(\bar{u}, \bar{v}, \bar{w}, \bar{p})$

## Usefulness

- If the averaged results are useful for the engineers
- i.e. the fluctuation are not interesting “only” their effect on the mean flow
- If modelling is accurate enough



# Eddy Viscosity modell

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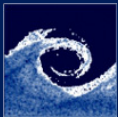
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### Idea

- Effect of turbulence is similar to effect of moving molecules in kinetic gas theory
- The exchange of momentum between layers of different momentum is by the perpendicularly moving molecules
- Viscous stress is computed by:  $\Phi_{ij} = 2\nu S_{ij}$





# Eddy Viscosity model (contd.)

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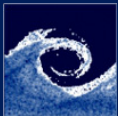
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## In equations...

- Only the deviatoric part is modelled
- The trace ( $k$ ) can be merged to the pressure (modified pressure), and does not need to be modelled
- Modified pressure is used in the pressure correction methods to satisfy continuity (see Poisson eq. for pressure)

$$\overline{u'_i u'_j} - \frac{2}{3} k \delta_{ij} = -2\nu_t \overline{S_{ij}} \quad (8)$$



# Eddy Viscosity

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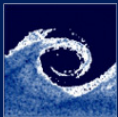
Viscosity is a product of a length scale ( $l'$ ) and a velocity fluctuation scale ( $u'$ )

- The length scale has to be proportional to the distance, what the fluid part moves by keeping its momentum
- The velocity fluctuation scale should be related to the velocity fluctuation caused by the motion of the fluid part

$$\nu_t \sim l' u' \quad (9)$$

## Newer results supporting the concept

Coherent structure view of turbulence, proves that there are fluid parts (vortices) which keep their properties for a while, when moving



# Two equations models

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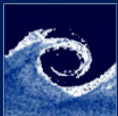
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- Length ( $l'$ ) and velocity fluctuation scales ( $u'$ ) are properties of the flow and not the fluid, they are changing spatially and temporally
- PDE's for describing evolutions are needed

### Requirements for the scales

- Has to be well defined
- Equation for its evolution has to be developed
- Has to be numerically 'nice'
- Should be measurable easily to make experimental validation possible



# k-e modell

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### Velocity fluctuation scale

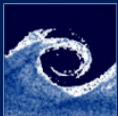
- TKE is characteristic for the velocity fluctuation
- It is isotropic (has no preferred direction)

$$u' \sim \sqrt{k} \quad (10)$$

### Length scale

- Integral length scale is well defined (see correlations)
- No direct equation is easy to develop
- Length scale is computed through the dissipation

$$\text{Recall: } \varepsilon = \frac{u_0^3}{l_0} \Rightarrow l' \sim \frac{k^{3/2}}{\varepsilon}$$



# Equation for the eddy viscosity

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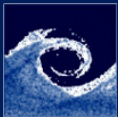
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$$\nu_t = C_\nu \frac{k^2}{\varepsilon} \quad (11)$$

$C_\nu$  is a constant to be determined by theory or experiments...

Our status...?

- We have two unknowns ( $k, \varepsilon$ ) instead of one ( $\nu_t$ )



# k model equation

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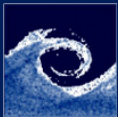
Equation for  $k$  was developed, but there are unknown terms:

$$\partial_t k + \overline{u_j} \partial_j k = \underbrace{-a_{ij} \overline{s_{ij}}}_{\text{Production}} + \underbrace{\partial_j \left[ \overline{u'_j \left( \frac{p'}{\rho} + k' \right)} - \nu \overline{u'_i s'_{ij}} \right]}_{\text{Transport}} - \underbrace{\varepsilon}_{\text{Dissipation}} \quad (12)$$

## Production

Production is directly computable, by using the eddy viscosity hypothesis

$$\mathcal{P} = -a_{ij} \overline{s_{ij}} = 2\nu_t \overline{S_{ij}} \overline{S_{ij}} \quad (13)$$



# k model equation

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## Dissipation

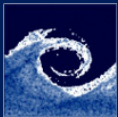
Separate equation will be derived

## Transport $\partial_j T_j$

- Can be approximated by gradient diffusion hypothesis

$$T_j = \frac{\nu_t}{\sigma_k} \partial_j k \quad (14)$$

- $\sigma_k$  is of Schmidt number type to rescale  $\nu_t$  to the required diffusion coeff.
  - To be determined experimentally



# Summarised k model equation

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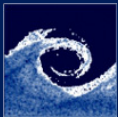
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$$\partial_t k + \overline{u_j} \partial_j k = 2\nu_t \overline{S_{ij}} \overline{S_{ij}} - \varepsilon - \partial_j \left( \frac{\nu_t}{\sigma_k} \partial_j k \right) \quad (15)$$

- Everything is directly computable (except  $\varepsilon$ )
- The LHS is the local and convective changes of  $k$ 
  - Convection is an important property of turbulence (it is appropriately treated by these means)





# Model equation for $\varepsilon$

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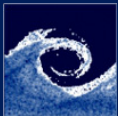
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- It is assumed that it is described by a transport equation
- Instead of derivation, based on the  $k$  equation

$$\partial_t \varepsilon + \overline{u_j} \partial_j \varepsilon = C_{1\varepsilon} \mathcal{P} \frac{\varepsilon}{k} - C_{2\varepsilon} \varepsilon \frac{\varepsilon}{k} - \partial_j \left( \frac{\nu_t}{\sigma_\varepsilon} \partial_j \varepsilon \right) \quad (16)$$

- Production and dissipation are rescaled ( $\frac{\varepsilon}{k}$ ) and 'improved' by constant coefficients ( $C_{1\varepsilon}$ ,  $C_{2\varepsilon}$ )
- Gradient diffusion for the transport using Schmidt number of  $\sigma_\varepsilon$
- **The  $\varepsilon$  equation is not very accurate! :)**



# Constants of the standard k-ε model

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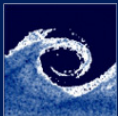
$$C_\nu = 0.09 \quad (17)$$

$$C_{1\varepsilon} = 1.44 \quad (18)$$

$$C_{2\varepsilon} = 1.92 \quad (19)$$

$$\sigma_k = 1 \quad (20)$$

$$\sigma_\varepsilon = 1.3 \quad (21)$$



# Example for the constants

## Homogeneous turbulence

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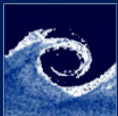
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$$d_t k = \mathcal{P} - \varepsilon \quad (22)$$

$$d_t \varepsilon = C_{1\varepsilon} \mathcal{P} \frac{\varepsilon}{k} - C_{2\varepsilon} \varepsilon \frac{\varepsilon}{k} \quad (23)$$



# Example for the constants

## Decaying turbulence

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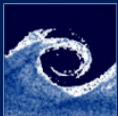
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Since  $\mathcal{P} = 0$  the system of equations can be solved easily:

- $k(t) = k_0 \left( \frac{t}{t_0} \right)^{-n}$
- $\varepsilon(t) = \varepsilon_0 \left( \frac{t}{t_0} \right)^{-n-1}$
- $n = \frac{1}{C_{2\varepsilon} - 1}$
- $n$  is measurable 'easily'



# $k$ - $\omega$ modell

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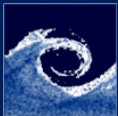
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- $k$  equation is the same
- $\omega \stackrel{\text{def}}{=} \frac{1}{C_\nu} \frac{\varepsilon}{k}$  Specific dissipation, turbulence frequency ( $\omega$ )
- equation for  $\omega$  similarly to  $\varepsilon$  equation
  - transport equation, with production, dissipation and transport on the RHS
- $\omega$  equation is better close to walls
- $\varepsilon$  equation is better at far-field

⇒ SST model blends the two type of length scale equations, depending on the wall distance



# Required Boundary Conditions

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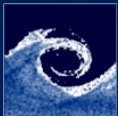
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The turbulence model PDE's are transport equations, similar to the energy equation

- Local change
- Convection
- Source terms
- Transport terms



# Inlet Boundary Conditions

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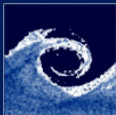
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- Neumann or Dirichlet or mixed type of BC can be used generally
- Inlet is usually Dirichlet (specified value)

### Final goal

- How to prescribe  $k$  and  $\varepsilon$  or  $\omega$  at inlet boundaries?



# Approximation of inlet BC's

## Turbulence intensity

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Inlet

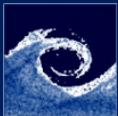
To use easy quantities, which can be guessed

Develop equations to compute  $k$  and  $\varepsilon$  or  $\omega$  from quantities, which can be guessed by engineers

Turbulence intensity

$$Tu \stackrel{\text{def}}{=} \frac{u'}{\bar{u}} = \frac{\sqrt{2/3k}}{\bar{u}}$$





# Approximation of inlet BC's

## Length scale

Turbulence  
II.

Miklós  
Balogh

Scales

TKE eq.

Modelling

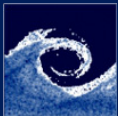
Boundaries

Inlet

### Length scale

$$l' \sim \frac{k^{3/2}}{\varepsilon} \Rightarrow \varepsilon$$

- From measurement (using Taylor hypothesis)
- Law of the wall (later)
- Guess from hydraulic diameter  $l \approx 0.07d_H$



# Importance of inlet BC's

## Turbulence II.

Miklós Balogh

Scales

TKE eq.

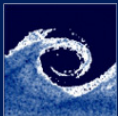
Modelling

Boundaries

Inlet

### If turbulence is governing a flow

- Example: Atmospheric flows, where geometry is very simple (flat land, hill) turbulence is complex
  - by spatial history of the flow
  - over rough surface
  - including buoyancy effects
- Sensitivity to turbulence at the inlet has to be checked
  - the uncertainty of the simulation can be recognised
  - measurement should be included
  - the simulation domain should be extended upstream



# Questions?

Turbulence  
II.

Miklós  
Balogh

Scales

TKE eq.

Modelling

Boundaries

Inlet

Thanks for your attention!